

# A Parallel Implementation of Three-Dimensional, Lagrangian, Shallow Water Equations

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Our work has focused on the  $f$ -plane shallow-water system:

$$h_t + (h\bar{u})_x + (h\bar{v})_y = 0. \quad (1)$$

$$h(\bar{u}_t + \bar{u}\bar{u}_x + \bar{v}\bar{u}_y) + \mathcal{C}^2 h(a_x + h_x) - fh\bar{v} = (\epsilon h(\bar{u}_x - \bar{v}_y))_x + (\epsilon h(\bar{u}_y + \bar{v}_x))_y - \frac{\epsilon k_1 h \bar{u}}{\lambda^2} + \frac{\epsilon k_2}{\lambda^2} (u_w - \bar{u}), \quad (2)$$

and,

$$h(\bar{v}_t + \bar{u}\bar{v}_x + \bar{v}\bar{v}_y) + \mathcal{C}^2 h(a_y + h_y) + fh\bar{u} = (\epsilon h(\bar{u}_y + \bar{v}_x))_x - (\epsilon h(\bar{u}_x - \bar{v}_y))_y - \frac{\epsilon k_1 h \bar{v}}{\lambda^2} + \frac{\epsilon k_2}{\lambda^2} (v_w - \bar{v}). \quad (3)$$

The fields  $(\bar{u}, \bar{v})(x, y, t)$  are the depth-averaged x-y components of the velocity field; the surface  $z = a(x, y)$  represents the bottom topography; the free surface of the fluid is given by  $z = a(x, y) + h(x, y, t)$  and the fluid occupies the region where  $h(x, y, t) > 0$ . We determine the free-boundary where  $h(x, y, t) = 0$  as part of the solution.

The turbulent viscous-stresses are given by

$$\Sigma = \epsilon h \begin{pmatrix} \bar{u}_x - \bar{v}_y, & \bar{u}_y + \bar{v}_x \\ \bar{u}_y + \bar{v}_x, & -(\bar{u}_x - \bar{v}_y) \end{pmatrix} \quad (4)$$

and this stress-tensor is purposely taken to be deviatoric and thus trace free. The terms  $(fh\bar{v}, -fh\bar{u})$  in (2) and (3) represent the Coriolis forces, the terms

$$\left( \frac{-\epsilon k_1 h \bar{u}}{\lambda^2}, \frac{-\epsilon k_1 h \bar{v}}{\lambda^2} \right)$$

represent the bottom friction forces, the terms

$$\left( \frac{-\epsilon k_2}{\lambda^2} (u_w - \bar{u}), \frac{-\epsilon k_2}{\lambda^2} (v_w - \bar{v}) \right)$$

represent the wind forces generated by wind velocities  $(u_w, v_w)$ . Finally  $\lambda^2 \ll 1$  represents the ratio  $(\frac{H}{L})^2$  where  $H$  is a typical vertical length scale and  $L$  is a typical lateral length scale and  $\mathcal{C}^2$  is the inverse Froude number.

Novel features of our work involve a ‘‘Lagrangian’’ reformulation of (1)-(3) and a robust Lagrangian code to integrate the system (1)-(3). We have developed all codes in MATLAB and have a parallel implementation of the code for the two-dimensional problem. These codes successfully reproduce the exact solutions obtained when the bottom surface is a paraboloid.